Physical processes associated with high primary production in Saanich Inlet, British Columbia

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Abstract: Saanich Inlet, British Columbia, has long been known for the presence in most years of anoxic bottom water. One factor contributing to this anoxia is a high level of primary production, which occurs as a major spring bloom followed by sporadic “mini-blooms” throughout the summer and early fall. The process(es) by which new production is re-fueled after nutrient exhaustion caused by the spring bloom are not well known, since Saanich is an inverse estuary and vertical mixing driven by winds and tides is low. This study presents new observational evidence that strongly suggests the dominant mechanism of nutrient resupply during the summer months is intermittent advective exchange, driven by pressure gradients set up by strong tidal mixing in passages outside Saanich Inlet itself. A simple box model is formulated to illustrate this mechanism. When driven by annual freshwater forcing and deepwater renewal functions characteristic of the region and measured tides for 1975, the model predicts resupply of nitrate during most of the periods observed in the 1975 observations of Takahashi et al. (1977). This “action-at-a distance” nutrient resupply mechanism, involving strong but localized turbulent mixing and subsequent distribution of the products of mixing over large horizontal distances by pressure-gradient-driven flow, is likely important in other coastal regions where estuarine circulation is weak.

1. Introduction

It has long been argued that primary production in Saanich Inlet is larger than that in typical B. C. coastal inlets, hence a major contributory factor to observed anoxia of bottom waters despite relatively regular deep-water renewal events (Anderson and Devol 1973). Timothy and Soon (2001) recently reported direct C\textsuperscript{14} uptake measurements of primary productivity contrasting Saanich Inlet with Jervis Inlet, a “normal” fjord-type estuary with oxygenated deep waters. Their multi-year observations confirm that average annual primary productivity in Saanich is indeed larger (roughly twice) than that in Jervis. An outstanding question is how such large primary productivity can be supported in an inlet which superficially might be expected to have rather low primary production.
In most fjord-type estuaries, a brackish surface layer flows outward as part of a normal estuarine circulation forced by dominant freshwater input at the head of the estuary. In such cases, biological new production is fueled by nitrate supplied by turbulent entrainment and pressure-gradient-driven upwelling of the base of the brackish layer. However Saanich Inlet is an inverse estuary, with freshwater supplied dominantly outside the inlet mouth, in winter by the near-field effect of the Cowichan...
River, in summer through the far-field effect of the massive freshet of the Fraser River (Figure 1). Thus mean surface-layer estuarine flow driven by freshwater forcing should be into rather than out of the inlet, hence associated with downwelling, not upwelling in the interior. In addition, because Saanich Inlet is not strongly forced by either winds or tides (Herlinveaux 1962), it has been assumed that the upper layers are characterized by low values of turbulent diffusivity. With inverse estuarine flow and low rates of turbulent nutrient resupply, high biological productivity in Saanich Inlet implies some other resupply process(es). Over time, the accumulation of various observations has suggested that possible mechanisms exist in the strong tidal mixing characteristic of regions just outside Saanich Inlet.

The first such evidence was provided by Takahashi et al.’s (1977) documentation of the detailed annual pattern of primary production in Saanich Inlet. Their study was the first to sample frequently enough (at least weekly) to describe not only the major bloom which occurs in spring and exhausts surface layer nitrate, but also a characteristic succession of subsequent “mini-blooms” which follow sporadic events of nitrate resupply during the summer and early fall. Noticing that increases in nitrate were correlated with increases in salinity (S), Takahashi et al. (1977) suggested that deeper, nutrient-rich waters were being moved nearer the surface by vertical mixing and/or vertical advection. Although no specific mechanism was suggested at the time, the time series observations of Takahashi et al. (1977: see their Fig.2) show a roughly fortnightly variability, suggesting some connection with the 14-day spring/neap tidal cycle. Indeed, Parsons et al. (1983) subsequently demonstrated significant anti-correlation between local tidal amplitude and the Takahashi et al. measurements of chlorophyll \(a\) at a period of 14 days. Parsons et al. (1983) went on to argue that the mouth of Saanich Inlet forms a frontal zone between stratified waters inside and strongly mixed (at least during spring tides) and nutrient-rich waters outside. The 14-day variability in primary production within Saanich could thus be driven by greater or lesser incursions of this front into the inlet as tidal velocities varied between spring and neap. This frontal mechanism almost surely contributes to the observed enhancement of primary production near the mouth relative to the head of Saanich Inlet (Timothy and Soon (2001) suggest a factor of 1.4 enhancement). However because maximum frontal excursion cannot exceed the
maximum tidal excursion (of order O(3 km)) in the vicinity of the mouth, this mechanism is unable to account for similar variation of nitrate resupply at stations near the head of Saanich Inlet, 20 km distant (see Sec. 4, Fig.7). Provided that primary production is reflected in export flux, the sediment trap measurements of Sancetta (1989) provide additional evidence that stations near the inlet head exhibit high frequency variability in primary production. This variability is likely associated with variable nitrate supply (keeping in mind that the 1 month sample interval of this study, considered normal (or even “high frequency”!) for such coastal studies, significantly aliases a 14-day period). Thus in addition to the localized resupply mechanism represented by frontal excursions near the mouth of Saanich Inlet, the observations require a mechanism which produces “action-at-a-distance”, i.e., a mechanism by which events outside Saanich can affect the entire inlet, not just areas nearest the mouth.

This purpose of this paper is to present new observational evidence (Sections 3 and 4) which clarifies the physical mechanisms involved in this hypothesized “action-at-a-distance”. While the suggested mechanism is physically quite different from that proposed by Parsons et al. (1983), the present study reinforces their general suggestion that tidally-modulated variability in stratification outside Saanich Inlet is a major determinant of the oceanography and the biological productivity within it. In Section 5, a simple box model, forced by observed tides and simplified freshwater fluxes, is used to illustrate the salient processes by which new nutrient is supplied to fuel the high primary production of Saanich Inlet.

2. Physical characteristics of Saanich Inlet and environs

Saanich Inlet is a 24 km long fjord-type inlet located at the southeastern end of Vancouver Island, British Columbia (Fig. 1a). While maximum depths exceed 200 m, a 75-80m deep sill at the mouth of the inlet restricts circulation of the deep basin waters (Fig.1b). Unlike most Pacific fjords, the waters of the deep basin are devoid of dissolved oxygen - anoxic - for most of the year. Saanich Inlet is unusual in another respect, namely, that physical processes and factors which force the circulation in most fjords are either absent or weak in this inlet.
Unlike normal estuaries, in which the major freshwater input occurs at the head of the inlet, there is relatively little freshwater discharge within Saanich Inlet itself because the catchment basin is rather small. The largest river flowing into the inlet (Shawnigan Creek, near the inlet mouth) has a maximum discharge of $O(4-5 \text{ m}^3 \text{ s}^{-1})$, predominantly in the winter rainy season. This local runoff is a minor percentage of two other sources of freshwater, both located outside the sill which defines the mouth of Saanich Inlet. The Cowichan River, entering 10 km to the northwest (Fig.1), has a winter maximum discharge of $O(20 \text{ m}^3 \text{ s}^{-1})$. Summer freshwater input to the region is dominated by the Fraser River, which discharges into the southern end of the Strait of Georgia within 50 km of the mouth of Saanich Inlet. Fraser River discharge contributes about half of the total freshwater supply to the Strait of Georgia (Griffin and LeBlond 1990), which peaks in a June freshet fuelled by melting snow pack in the interior of the province. Maximum summer Fraser discharge of $O(8000 \text{ m}^3 \text{ s}^{-1})$ vastly exceeds that of even maximum (winter) values of the Cowichan, so that the Fraser River “far-field” source dominates summertime freshwater forcing of Saanich Inlet. The freshet “pulse” moves seaward from the southern Strait of Georgia to Juan de Fuca Strait and eventually out to the Pacific Ocean through a complex region of tidal channels among the Canadian Gulf Islands and the American San Juan Islands. Saanich Inlet is connected to Boundary Passage and Haro Strait, the major inter-island passages through this region, by Satellite and Swanson Channels (Fig.1).

Both wind and tidal forcing of turbulence within Saanich Inlet are generally weak. Because the north/south axis of the inlet is roughly normal to prevailing winds, the wind field is much less energetic in Saanich than in most other B.C. inlets, which are oriented more nearly east/west. There are no prominent diurnal sea-breezes during summer months, nor is the inlet subject to strong katabatic outflows of cold arctic air during winter months. Tidal currents are weak throughout Saanich Inlet. Moreover since the sill is moderately deep and the mouth of the inlet rather wide (3.5 km), the energetic interactions of a stratified flow with topographic constrictions (Farmer and Freeland 1983) which are common in many other Pacific fjords are not observed in this inlet. The weakness of wind and tidal forcings is consistent with available direct observations of turbulence properties. For several years, Saanich Inlet acted as a pre-cruise test site for
microscale shear sampling which provides profiles of turbulent kinetic energy dissipation rate $\varepsilon$, a commonly used metric for the strength of turbulent mixing (Tennekes and Lumley 1972). As seen in the typical profiles of Fig. 2a, observed values of $\varepsilon$ rarely rise above a system noise level of $O(10^{-9} \text{ m}^2\text{s}^{-3})$ which translates, through standard microscale estimation (Gargett 1997), to a diffusivity of $O(2\times10^{-6} \text{ m}^2\text{s}^{-1})$. Using a mean nitrate gradient of $O(0.5\text{ mmol m}^{-2}) = 10\text{ (mmol m}^{-3})/20\text{ m})$ typical of mid-inlet (see Fig. 6, Stn. S5), this diffusivity would produce a vertical nitrate flux of only $O(10^{-1} \text{ mmol m}^{-2}\text{d}^{-1})$, much too small to account for measured fluxes of $O(4 \text{ mmol m}^{-2}\text{d}^{-1})$ during at least part of the fortnightly tidal cycle (see Section 4). Moreover, turbulent nitrate flux would be expected to be particularly weak during the summer months, due to a combination of weak winds and strong near-surface stratification caused by heating as well as freshening of the surface layer.

![Figure 2](image)

Figure 2: Typical depth profiles of the logarithm of the turbulent kinetic energy dissipation rate $\varepsilon$, as determined from airfoil probes carried on a microscale profiler in (a) Saanich Inlet (Station S5) and in (b) the Boundary Pass/Haro Strait tidal fronts that occur near the entrance (Station S9) to Swanson Channel, the major “deepwater” pathway to Saanich Inlet. Assuming quasi-steady-state turbulence, the 3-4 orders of magnitude difference in $\varepsilon$ between these two regimes implies similar difference in the rates at which energy is being supplied to turbulent motions.

While turbulence is weak within the inlet, it is anything but weak in many adjacent channels and passages. A large tidal volume moving through contorted passages and over sills produces strong (up to 4 knots) tidal currents that generate intense mixing in localized areas. Nearby locations (see Fig.1) where strong mixing has been
documented include Sansum Narrows (Gargett 1994), Cordova Channel (Lu and Lueck 1999), and recurring tidal fronts (Gargett 1994, 1999: Farmer et al. 1994) in the region of Boundary Pass/Haro Strait near the mouth of Swanson Channel (see Fig.1). As illustrated in Fig.2b, turbulence in such areas can reach dissipation levels that are 3-4 orders of magnitude above the noise levels observed within Saanich. This tidally-generated turbulence varies with the strength of tidal flows, hence is modulated at all significant tidal periods. At time scales longer than the dominant diurnal (K1) and semi-diurnal (M2) periods, tidal flows vary strongly with the fortnightly cycle that gives strongest flows during spring tides, weakest flows during neap tides. Measurements of $\epsilon$ made in Cordova Channel by Lu and Lueck (1999) provide direct evidence for an associated spring/neap cycle in tidal mixing. At longer periods, tidal mixing is expected to vary at the semi-annual period which gives the strongest tidal flows of the year in June and December.

Because tidal mixing works against the stabilizing effect of seasonal freshwater input (and additional surface heating in the summer months), tidally driven variations in turbulence intensity result in large variations in density stratification in the region. Direct evidence of such variation is shown in Fig.3. The color-coded background shows the evolution of the field of $\log \epsilon$ (determined by an acoustic large-eddy method: Gargett 1999) across a strong tidal front in Haro Strait over a four day period as the tides increase from neap towards spring. Superimposed vertical profiles of density ($\sigma_i$) show that the late winter density structure, quasi-two-layer at the time of the neap tide, evolved into a much weaker, nearly linear stratification after only four days of increasing tidal flows and mixing. On maximum spring tides, water column density may be nearly homogenized within such localized turbulent structures. Griffin and LeBlond (1990) provide indirect evidence that variation of turbulent mixing at the spring/neap periodicity strongly modulates seaward freshwater export in the upper layers of Juan de Fuca Strait.
Figure 3: Colour codes the field of $\log \varepsilon$, determined from a new acoustic large-eddy technique (Gargett 1999), measured during approximately north/south (left to right) transects of a strong tidal front in Haro Strait as tides increased from neap towards spring. Superimposed vertical profiles of $\sigma_t$ (courtesy of J. Moum, Oregon State University) show the effects of tidal mixing on the late winter density structure, quasi-two-layer at the beginning but evolving quickly to a much weaker, nearly linear stratification after only four days of increasing tidal flows and turbulent mixing.

3. Observational evidence of spring/neap tidal forcing mechanism of mean flows

In July 1994, Canadian government agencies together with organizations of concerned citizens began the Saanich Inlet Study (SIS), a broadly-based examination of the present state of the marine environment of the inlet and its ability to assimilate
wastes from future development around its shores. During the SIS, a short-term measurement program was undertaken to address gaps in oceanographic data relating to circulation and nutrient transport above sill depth. The results presented here came from an intensive sampling period from July 5 to 20, 1995. Both shallow-drogued drifter tracks and water property variations measured during this period provide clear evidence of the dominant process that drives the shallow circulation.

Figure 4: Time histories of radio-tracked surface drifters, drogued at 2 m to reveal surface layer motion, over 5 day intervals during (a) neap and (b) spring tides.

Over the period July 10th to 20th, sets of drifters drogued to follow the top 2m of the water column were deployed and tracked. During the week of neap tides (Fig.4a), drifters deployed in Finlayson Arm, the narrowest part of the inlet, exhibit the back and forth motion associated with the semi-diurnal tide, superimposed on a slight tendency for net up-inlet (southerly) movement. At the broader northern end of the inlet, most of the drifters appear locked in a CCW rotating “eddy” associated with weak inflow(outflow) on the western(eastern) side of the inlet. No drifters left the inlet, although several grounded on the eastern shore.
In startling contrast (Fig. 4b), most of the drifters deployed in the northern part of the inlet as tides built towards the spring tidal maximum moved swiftly northward and out of the inlet. Of the drifters deployed in the lower reaches of the inlet, most grounded immediately. However one drifter set in Finlayson Arm moved back and forth for a couple of days and then proceeded northward the entire length of the inlet. During this period of time, winds were generally weak, hence could not be responsible for the observed northward transport of floats.

The combined drifter results suggest that spring/neap variation in tidal flows somehow produces strong modification of non-tidal horizontal flow, resulting in periodic reversals of the weak surface inflow which characterizes an inverse estuary like Saanich. Moreover such reversals involve the entire inlet, not just the regions nearest the mouth. Reversal periods, during which Saanich Inlet temporarily behaves as a normal estuary, will be associated with a “normal” estuarine upwelling/entrainment of nutrient-rich deep waters which flow in at depth to compensate for the surface-layer outflow.

The mechanism involved in translating tidal variability into mean flows is suggested by salinity sections that were taken from the head of Saanich Inlet out through the connecting passages (Satellite and Swanson Channels) to the major external deepwater channel (Haro Strait/Boundary Pass). Figs. 5a and 5c show sections taken during two neap tides bracketing the spring tide section of Fig. 5b. The first neap tide survey (July 6th) shows waters with S < 29.0 occupying the top 20 m of most of the inlet. One week later, at the time of spring tides, only a vestige of waters with S < 29.0 remained in the inlet. The layer of salinities between 29.0 and 30.0 had thickened within the inlet, as the 29.0 isohaline moved upward in the water column, and disappeared entirely outside the inlet. The following neap tide survey (July 20) shows that water with S < 29.0 had reoccupied the entire length of the inlet. Local runoff cannot account for the observed replenishment of the low salinity layer during the second neap tide period. Only 10.3 mm of rain fell in the area from the last week in June to the last week in July, with most falling on July 9th and 10th, just before the spring-tide flushing of low salinity water from the inlet. In addition, other water properties unaffected by precipitation
(temperature and dissolved oxygen, unpublished data) show similar large scale changes in the three surveys.

Figure 5: Sections of salinity $S$ taken from the head of Saanich Inlet (left) through the extended sill region outside its mouth and into deeper water near the junction of Swanson Channel and Haro Strait (right). Two sections taken during neap tides (a) July 6th and (c) July 20th are separated by a section taken during the intervening spring tide on (b) July 13th. Note the virtual disappearance of low salinity water ($S < 29.0$) from Saanich Inlet and the weakness of salinity stratification outside the inlet in the spring tide section.

Because density is determined predominantly by salinity in this area, there are major changes in stratification associated with the observed changes in the salinity field. While quantitative details differ, both neap tide sections reveal strong near-surface density stratification through the entire section. On spring tides, there is a reduction of stratification within the inlet, but the largest changes are observed in the density gradients outside the inlet. Stratification is greatly reduced both in the shallow approaches (Satellite Channel) and the deeper connection (Swanson Channel) to Haro.
Strait (far righthand side of Fig.5), where neap tide salinity differences of ~ 2.5 – 3.0 from surface to 100m become differences of only ~ 0.5 at spring tides. In the face of continuing freshwater input from the Fraser River, the observed spring tide reduction in stratification in these areas external to Saanich Inlet can only result from increased tidal mixing, as documented in the previous section.

4. Time series of nitrogen species in Saanich Inlet

To this point, data and discussions have dealt with the physical circulation and water property changes which occur inside Saanich Inlet in response to changing density and stratification outside. We now turn our attention to the effects of this variation in the physical system on transport and cycling of nitrogen and the biological productivity of this fjord.

Because nitrate supply periodically limits summer phytoplankton growth in Saanich Inlet (Takahashi et al. 1977), it is necessary to understand the mechanisms by which nutrients are resupplied to the euphotic zone before we can describe plankton dynamics over the entire summer growth season. Consequently, the SIS included a sampling program aimed at resolving the transport of nitrogen in its biologically active forms over the spring/neap cycle. Three stations (S3, S5 and S6, see Fig.1) were sampled every two to three days between July 5th and 20th. Discrete water samples were analyzed for nitrate (NO$_3$) and ammonium (NH$_4$) (Barwell-Clarke and Whitney 1996), and for dissolved organic nitrate (DON) (Solorzano and Sharp 1980). Sea water samples from S3 and S6 were filtered to determine concentrations of particulate organic carbon and nitrogen (POC/PON) using a CEC440 elemental analyser. In addition to the water sampling program, a sediment trap moored near S3 at a depth of 50 m was sampled every two or three days to determine vertical transport of particulate nitrogen. Measurements show that of the available sources, NO$_3$ and DON together, in approximately equal proportions, account for 80-90% of biologically active nitrogen. Moreover NO$_3$ and DON co-vary temporally; we thus use NO$_3$ concentration alone in the following discussion of the temporal variability of biogenic nitrogen at the three stations.
Figure 6: Top panel: tidal height in Saanich Inlet, showing the strong neap/spring variation during the period of measurements. The gray-coded panels are time series of NO$_3$ concentration at three Saanich stations (see Fig.1 for locations). Each station shows a pronounced neap/spring variation in near-surface NO$_3$.

As seen in the gray-scale panels of Fig.6, nitrate concentrations exhibit the same overall temporal pattern at all three stations, with quantitative differences related to the distance of each station from the mouth of the inlet. At all stations, nitrate concentrations in the surface layer (0 to ~ 20 m) increase during spring tides (top panel). This increase is largest in magnitude and earliest in time at the sill station (S6),
appearing at interior stations progressively later and with decreased magnitude. However even at the southernmost interior station (S3), there is still distinct spring-tide enrichment of surface nitrate.

Fig.7 shows that the increase in surface-layer nitrogen at S3 is highly correlated with salinity S (which is in turn highly anti-correlated with temperature T during the summer period of these measurements), as would be expected if the surface layer were advecting out of the inlet, and being replaced by upwelled deeper waters. The effect of this resupply of nitrate to the surface layer may be seen in the dramatic increases in both the concentration of suspended particulate organic nitrogen (PON: Fig.7b) and the rate of PON export, as measured at the 50m sediment trap (Fig.7c), that are observed at the end of the time series, several days past peak spring tides. Such a time lag between resupply of a limiting nutrient and increase in suspended particulates (phytoplankton) is to be expected, given finite times required for phytoplankton growth. Export production has a potentially slower timescale, associated with processing of the subsequent burst of primary production by the resident zooplankton. However the rise in PON and that in export production are observed to occur nearly simultaneously, at least as well as can be determined from the time resolution of the water and trap sampling sequence. Absence of a significant time lag between these two variables suggests that most export production is associated with rapidly sinking fecal pellets produced by macrozooplankton, a conclusion consistent with analyses of trap material collected in this study and in former studies in Saanich Inlet (Sancetta 1989).
Figure 7: Time series of (a) temperature $T$ and salinity $S$, (b) nitrate $\text{NO}_3$ and particulate organic nitrogen $\text{PON}$ averaged over the upper 35 m of the water column from samples collected at 0,5,10,15,20,30 and 40 m and (c) sedimentation rate of $\text{PON}$ measured by a sediment trap moored at 50m depth. All measurements were made at Station S3.

The above interpretation assumes that the observed changes in particulate and export nitrogen are produced by local phytoplankton growth and subsequent zooplankton grazing, fueled by vertical resupply of $\text{NO}_3$, a limiting nutrient. However because the time lag observed between the increase in $\text{NO}_3$ and subsequent increases
in particulates happens to coincide roughly with the time interval between spring and neap tides, an alternate interpretation could be that the return to an inverse estuarine circulation characteristic of neap tides has advected increased particulates to the trap site from further north in the inlet. We think this advective explanation is unlikely, given the drifter evidence of weak recirculating flows during neap tides (Fig. 4a).

5. Box model of Saanich Inlet and environs

The observed changes in large-scale salinity distributions, reinforced by the surface drifter measurements, make it clear that there are large scale movements of low salinity surface waters out of the inlet, then back into it, over the course of a two week spring/neap tidal cycle. Observed changes in biogenic nitrogen suggest that these motions are associated with major impacts on the summertime biological productivity of Saanich Inlet. The fact that such impacts are observed at station S3, well away from the mouth of the inlet, indicates that the process that generates the horizontal motions must act over the entire inlet, not merely in frontal regions near its mouth. We suggest that pressure forces resulting from tidally-forced changes of stratification outside the inlet are responsible for generating surface-layer outflow during spring tides. Recalling that the dominant source of summer freshwater is outside Saanich Inlet, surface layer salinities outside the inlet fall below those inside the inlet (at the same depth) when tidal mixing is weak during neap tides. The resulting horizontal density (salinity) difference provides a pressure force that causes flow into the inlet in a surface layer and out of it at depth, i.e., the inverse estuarine circulation typical of Saanich Inlet. However as tidal flows increase towards spring tides, strengthened vertical mixing outside the inlet erodes surface layer salinities outside the inlet, but not within it. If enhanced tidal mixing causes vertical salinity fluxes that are sufficiently large compared to the continuing freshwater flux, surface-layer salinities outside the inlet may fall below those within it, reversing the pressure gradient and hence the flow directions. This pressure-gradient mechanism accounts for the major observational features described in the previous sections, namely the inlet-scale response and the association of flow reversals with spring-tide induced changes in stratification outside the inlet.
A simple box model has been formulated to illustrate the interaction of stabilizing freshwater forcing and destabilizing tidal mixing in the operation of this mechanism. The model is roughly scaled to the dimensions of Saanich Inlet and used to predict periods of outflow, hence nutrient resupply, during the summer of 1975, when the nutrient measurements of Takahashi et al (1977) were taken. The box model, shown schematically in Fig.8, consists of two ‘basins’, each represented by a thin upper layer, extending from the surface to 20m, approximately the depth of the halocline within Saanich Inlet, and a lower layer of thickness $h = 60$ m extending from 20 m to 80 m, an approximate depth of the channels outside its mouth (labeled generically ‘Haro Strait’). The inlet is assumed to be 20 km long ($L$) and an average 2 km wide. The adjoining region represents neighboring areas such as Satellite and Swanson Channels as well as parts of Haro Strait itself. In the model results presented here, the effective area of this region is taken equal to that of Saanich, ie $\Delta = 1$.

![Figure 8: Schematic of box model of Saanich Inlet and environs: for details see text and Appendix.](image)

Mass transport $q \equiv Q_o (\rho_2 - \rho_1)$ is assumed to vary with a pressure head associated with the difference in densities of the two upper boxes. This parameterization, due originally to Stommel (1961), has been widely used in studies of
the thermohaline circulation (Whitehead 1995), and has recently been applied to shallow inverse estuaries (Hearn 1998). The transport $q$ is defined as positive for outward surface layer flow, as illustrated in Fig. 8. The horizontal flow between the surface boxes is assumed to be accommodated by up/downwelling across vertical box boundaries and return horizontal flow between the lower boxes. Thus, with surface layer transport in the direction shown, water is advected into the deep box (3) of Saanich Inlet and upwelled into the surface box (1).

Density variations are assumed to be dominated by variations in salinity (although temperature ($T$) also influences density, atmospheric forcing of $T$ is uniform over the limited horizontal scale of this model, so will not contribute to differential forcing of surface layer densities). Since the model is focussed on the summer phytoplankton growth season, (local) winter freshwater inputs are ignored, and the model is forced by equivalent salt fluxes associated with (a) the annual freshet of the Fraser River, input to the upper box of the outer basin, and (b) a late summer input of saline water to the lower outer box. The latter salt flux is associated with the annual appearance in Haro Strait of saline California Current waters that upwell onto the outer shelf of Vancouver Island in the summer and are entrained into the estuarine inflow that proceeds landward in the deep Strait of Juan de Fuca (Thomson 1981). In the normalized forcing functions, shown in Fig. 9a, the amplitude of the saline inflow function is set equal to that of the freshwater function in order to preserve annually-averaged mean salinity in the system (for details see Appendix). However temporal variation occurs within the annual cycle because of the time lag $\tau$ between the two processes. In the simulation shown, freshet starts at the end of May and peaks at the end of June; $\tau$ is taken to be two months.

Further details of the full equations for conservation of salt in this box model system, as well as the salt flux forcing functions used to drive it, are found in the Appendix.

Finally, it is necessary to parameterize $K_v$, the vertical diffusivity of tidal mixing in the model. Between the two inside boxes, we set $K_v = 0$, consistent with previously discussed evidence for low turbulent diffusion within Saanich Inlet. Between the two outside boxes, we use the relation of Osborn (1980)
relating turbulent kinetic energy dissipation rate $\varepsilon$ and the buoyancy frequency $N = (-g \rho_o^{-1} \partial \rho / \partial z)^{1/2}$. Defining $u_e$ as a characteristic speed and $t_e \sim u_e / l_e$ as a characteristic time scale for the large eddies (length scale $l_e$) of the turbulent field and substituting the standard scaling relation $\varepsilon \sim u_e^2 / t_e$ (Tennekes and Lumley 1972), (1) becomes

$$K_v \sim \frac{\varepsilon}{N^2}$$

(1)

an expression that is closed by assuming that $u_e \propto U$, where $U$ is a typical tidal speed, and that $t_e \propto N^{-1}$, ie that the turbulent overturning time scale is set by the local buoyancy period. The latter assumption implies continual strong buoyancy control of turbulent motions, since it implies that a turbulent Froude number $Fr_e = u_e / N l_e$ is always of $O(1)$.

An annual time series of $U(t)$ is determined by first-differencing observed tidal heights at a tidal station (Fulford Harbour) just off Satellite Channel. Then, using $N \sim (\Delta \rho)^{1/2}$ where $\Delta \rho = (\rho_4 - \rho_2)$ is the density difference between the surface and deep boxes in ‘Haro Strait’,

$$K_v \sim \frac{U^2}{N} \quad \text{or} \quad K_v = K_o \frac{U^2}{(\Delta \rho)^{1/2}}$$

(3)

where $K_o$ is a coefficient necessary to give magnitude (and the appropriate dimensions) to the scale relationship in (3). This expression for $K_v$ involves system properties that are both external ($U$) and inherent ($\Delta \rho$), in the sense of evolving with evolution of model properties. Imposing such (computed) vertical diffusion at all times was found to produce very smooth evolution of box salinity fields (not shown), unlike the on/off behavior suggested by the observations. We obtained more abruptly varying flows, consistent with the character of the observations, by adding an *ad hoc* threshold
criterion\textsuperscript{1}, requiring that the diffusivity given by (3) be imposed only when a critical Reynolds number $Re_c$ is exceeded by a mean-flow Reynolds number $Re = U H / \nu$ based on the typical barotropic tidal speed $U$ used in (3) and the depth $H = h(1 + \delta)$ corresponding to minimum water column depth in the approach passages.

The adjustable parameters of the model are $Q_o$, the coefficient that sets the magnitude of flow resulting from a given pressure (density) difference, the magnitude $F_o$ of the freshet forcing function, the diffusion coefficient $K_o$, and a value for the critical Reynolds number $Re_c$. We first estimated $F_o$ using an average value ($Q_r = 6000 \text{ m}^3\text{s}^{-1}$) of the Fraser River peak freshet flow, transformed to an equivalent salt flux through the relationship $F_o = Q_r S_o / V_o$ (Gill 1982: see Appendix). We then chose values for $Q_o$ and $K_o$ which produced observationally reasonable summer salinities in the Saanich Inlet boxes for which we have the most observational information. Although a full discussion of parameter sensitivities of the box model is beyond the scope of this paper, it should be noted that were the model forced with larger or smaller values of Fraser flow, results similar to those presented here can be achieved by adjustments to $Q_o$ and/or $K_o$.

Alternately leaving the chosen parameters constant while allowing $F_o$ to vary over the range observed for peak Fraser River freshet flow, model results for the box salinities vary within observationally reasonable ranges. The final parameter, $Re_c$, has relatively little effect on the annual cycle of mean salinities, but more substantial effect on the difference $(S_2 - S_1)$ between the salinities of the upper layers, hence on the potential...

\textsuperscript{1} It is more usually assumed that the appropriate condition for initiation of turbulence in a stratified fluid is that of $Fr = 1$. However a preliminary attempt to “turn on” turbulent mixing with a constant (large) diffusivity only when a Froude number criterion was met led to a phenomenon labelled “runaway stratification” by Simpson et al. (1990). After an initial period during which the stratification varied (a period whose length depends upon the initial imposed stratification), the system eventually achieved a degree of stratification which, given finite tidal velocities and a constant diffusivity, could no longer be reduced sufficiently to trigger turbulent mixing. After this point, the stratification increased monotonically, rapidly reaching totally unrealistic values. In contrast, our parameterization implies that a sufficiently high $Re$ will initiate turbulence regardless of stratification. Once initiated however, the resulting turbulence is strongly influenced by the stratification, through the requirement that the turbulent Froude number $Fr$ is always of $O(1)$.

We offer no justification for this very different parameterization beyond the obvious - it works (in the sense of providing stable model solutions), while parameterization based on $Fr$ doesn’t.
for periodic flow reversals. With $Re_c = 0$, which corresponds to $K_v$ “on” at all times, the flow is into the inlet at the surface from April until mid-August, when it reverses to a (winter) condition of surface outflow, ie flow direction reverses only once a year. Increasing $Re_c$ gradually, we find that values of $Re_c$ around 0.5 produced occasional flow reversals during the summer season. As $Re_c$ approaches 1 (corresponding to $K_v = 0$ at all times), the system approaches an uninteresting advective limit.

As an example of the behavior of this model, Fig.9 shows results forced with the annual time series of the freshwater and deepwater forcing functions shown, and values of $F_o = 5.3x10^{-6} [S] s^{-1}$, $Q_o = 2x10^{-3} m^3 s^{-1}$, $K_o = 1x10^{-3} (m^{7/2}s^{-2}kg^{-1/2})$, $Re_c = 0.4$. To compare at least qualitatively with observations taken by Takahashi et al.(1977) in 1975, tidal heights observed during 1975 are used to determine $U$, as described above. Fig.9b shows the resultant time series of $Re/Re_c = U/U_c$: turbulent diffusion with $K_v$ calculated from eqn.(3) is turned on only when $Re/Re_c > 1$ (ie above the horizontal line in Fig.9b). The resultant annual cycle of salinities in the four boxes, calculated from an initial state in which both basins are weakly stratified, is seen in Fig.9c. Fig.9d shows the associated evolution of stratification, represented respectively in “Haro Strait” and “Saanich Inlet” by the vertical salinity differences $(S_4 - S_2)$ and $(S_3 - S_1)$, as well as the sense of advective flow $q = Q_o(\rho_2 - \rho_1) \propto (S_2 - S_1)$ forced by the horizontal density (salinity) difference between the two upper boxes.

Arrival of the Fraser River freshet is signaled first by a strong decline in the salinity $(S_2)$ of the upper Haro box, quickly communicated to the upper Saanich box $(S_1)$ by advection in the inverse-estuary sense. During this period of early June to late July, stability increases in both basins: however there are brief intervals (shaded in Fig.9d), coincident with spring tides, when tidal mixing in the Haro boxes is sufficiently strong that $S_2 > S_1$ and the advective flow temporarily reverses, ie the surface layer flows out of Saanich Inlet. During late July and August, as the freshet wanes but before the increase in stratification outside the inlet associated with arrival of the saline deepwater pulse, spring tides of similar or even lower magnitude lead to increasing periods of surface layer outflows. With the arrival of the saline deepwater pulse, $S_4$ increases rapidly. As this salt is mixed upwards on successive spring tides, increasing the density
Figure 9: (a) Equivalent salt forcing functions (F) normalized by $F_0$, an amplitude derived from Fraser River freshet peak flow: freshwater input appears as negative salt forcing. (b) The function $Re/Re_c$ used to turn 'Haro Strait' mixing on/off (for details see text). (c) Predicted annual cycles of salinities in the four model boxes. (d) Predicted annual cycle of density stratification inside and outside Saanich Inlet, represented respectively by salinity differences ($S_3 - S_1$) and ($S_4 - S_2$). Surface layer flow is out of Saanich Inlet when $(S_2 - S_1) > 0$.

of the surface “Haro Strait” box, the outflow periods in which $S_2$ exceeds $S_1$ lengthen, and finally become continuous by the mid-winter (December) semi-annual maximum of the tides.

As seen in Fig.9, this model predicts summer outflows, of increasing duration, associated with spring tides in late June/early July, early August and early September, roughly coincident with all but one of the sharp increases of nitrate observed during the
summer measurements of Takahashi et al. (1977). Closer comparison of model predictions with measurements is not warranted, because deepwater renewal, one of the two major physical forcing functions for the model, is poorly constrained by observations. While the timing of potential renewal events is set by the tidal velocities, which vary relatively little from year to year, the occurrence and/or duration of actual renewal events during a particular year depends upon magnitude and timing of both the Fraser freshet and the deepwater renewal event. While characteristics of the Fraser freshet are well-determined by river flow monitoring dating back to 1912, there are no corresponding records of timing or magnitude of the dense deepwater inflows. All that is known about these inflows is that they generally occur in late summer or early fall, and must on average balance the net summer annual freshwater inputs, since the mean salinity of Saanich Inlet has not been observed to change over the extent of historical measurements. Further work in progress will attempt to determine which features of the physical forcing functions produce the most variation in timing and duration of the sporadic summer nutrient renewal events that make a major contribution to the summer productivity of this inlet.

6. A sedimentary record of summer mini-blooms?

The anoxic bottom waters of Saanich Inlet prevent the usual macro-fauna from colonizing its sediments, which thus remain undisturbed by burrowing and digging animals. Gross et al. (1963) described the resulting seasonal strata preserved in these sediments, seen as a regular sequence of dark and light bands (varves) in sediment cores. An annual cycle consists of one light and one dark band, resulting respectively from deposition of biogenic material (light) produced during the summer growing season and of terrigenous material (dark) washed into the inlet by winter runoff. Gross et al. (1963) also remarked that the light-colored biogenic (summer) laminae commonly contained even finer layers. Fig.10 contains examples of such micro-varves, seen in
Figure 10: X-radiographs over several of the annual varves seen in replicate freeze-cores taken at a Saanich Inlet site roughly midway between Stations S4 and S5 (Fig. 1). Very thin micro-varves can be distinguished within most of the light (summer) bands: these may result from quasi-periodic "mini-blooms" resupplied with nutrients via the physical processes described in this paper. (Core radiographs courtesy of V. Tunnicliffe, University of Victoria).

multi-year segments of modern freeze-cores, while other examples may be found in McQuoid and Hobson (1997) and Dean et al. (2001). Viewed as individual light layers superimposed upon a dark background, this fine detail within the productive summer season could be sedimentary evidence of summer "mini-blooms" produced by semi-
periodic nutrient resupply by the tidally-forced renewal process discussed in this article. Viewed alternately as individual dark layers superimposed upon a light background, the micro-varves might equally well be evidence of the effects of intense local run-off events superimposed upon biological production that is relatively constant throughout the summer growing season. The latter interpretation seems much the less likely, both because of the existing observational evidence of episodic biological production, and because Saanich Inlet lies within a region with normally dry summers (Herlinveaux 1962). However future work will attempt to differentiate between these two possibilities by examining historical records of both Fraser River flow and local rainfall in relation to the sedimentary record of microvarves.

7. Discussion and conclusions

We feel confident that the observations and model presented here provide a major part of the explanation for the high primary production known to characterize Saanich Inlet. Enhanced primary productivity results from a combination of effects. First, the sheltered nature of the inlet with respect to summertime winds and tidal mixing leads to a stably stratified environment in which phytoplankton are retained in a shallow, well-lighted surface layer. Next, surface-layer macronutrients depleted by phytoplankton blooms triggered by these favourable light conditions are resupplied on a semi-regular basis throughout the growing season. The resupply mechanism is inlet-scale upwelling, driven by pressure-gradient forces set up by effects of strong spring-tidal mixing on vertical density gradients outside the inlet. Legendre (1981) argued that regions of especially high primary productivity must be characterized by a temporal sequence of locally intense mixing followed by stable stratification. In Saanich Inlet, the separation between favourable nutrient and light conditions is spatial rather than temporal in nature, with advection an essential connecting element. Finally we note that in the high-growth periods between nutrient renewal events, weak inverse estuarine circulation will retain local primary production within Saanich Inlet, as well as import material from the frontal region at its mouth (Parsons et al. 1983, Hobson and McQuoid 2001). Coupled with high primary productivity, this retention mechanism also contributes
to the high levels of secondary and export production which are observed in Saanich Inlet.

The present study of Saanich Inlet documents a system in which nutrients necessary for biological production are supplied to the euphotic zone of a highly productive region through processes that act outside the region itself. The process by which intense but highly localized mixing affects phytoplankton production over large spatial scales through horizontal redistribution in an estuarine circulation is not unique to Saanich Inlet. Although not concerned with biological consequences, Stacey and Gratton (2001) have recently modeled “reverse renewal” in the two-silled fjord of the Saguenay, off the St. Lawrence estuary, as a result of pressure gradients associated with more vigorous $M_2$ tidal mixing in the outer basin than the inner basin. In another example, Crawford (1991) estimated that the dominant supply of nutrients to the highly productive ecosystem on the southern British Columbia shelf is through advection in the surface layer that exits the Strait of Juan de Fuca at the south end of Vancouver Island (Fig.1). Nutrients are mixed into this layer predominantly within the inter-island passes and channels at the southeast end of the Strait of Georgia, far from the outer shelf region where they are utilized. We suggest that the “action-at-a-distance” mechanism of nutrient resupply to phytoplankton populations that we have documented in Saanich Inlet is likely typical of estuarine and coastal regions of strong freshwater influence.

Finally, we warn that marine ecosystems in estuarine areas that are strongly influenced by turbulent mixing can be severely impacted by potential effects of global warming (increased surface temperatures, increased or decreased precipitation/run-off) on the stratification of the environment in which they are embedded.

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Appendix:

The 4-box model configuration is shown in Fig. 8. For Box j, we denote thickness as $h_j$ and area as $A_j$. Box salinity (salt/unit mass) is given by $S_j$ and the mass of water in the jth box is $m_j = \rho_o V_j = \rho_o h_j A_j$, where $\rho_o$ is a reference density. Equations for the time rate of change of the mass of salt $m_j S_j$ in the jth box are given by

$$m_j \frac{dS_j}{dt} = q_o^+(S_3 - S_1) + q_o^-(S_2 - S_1) + \rho_o w_i A_i (S_3 - S_1) \quad (A1)$$

$$m_2 \frac{dS_2}{dt} = q_o^+(S_1 - S_2) + q_o^-(S_4 - S_2) + \rho_o w_o A_o (S_4 - S_2) - m_2 F S_2 \quad (A2)$$

$$m_3 \frac{dS_3}{dt} = q_o^+(S_4 - S_2) + q_o^-(S_1 - S_3) - \rho_o w_i A_i (S_3 - S_1) \quad (A3)$$

$$m_4 \frac{dS_4}{dt} = q_o^+(S_2 - S_4) + q_o^-(S_3 - S_4) + \rho_o w_o A_o (S_4 - S_2) + m_4 F S_4 \quad (A4)$$

Advection between boxes is characterized by a flux of mass $q_o = Q_o (\rho_2 - \rho_1)$ proportional to the density difference between the two upper layers (Stommel 1961). To allow for reversal of the flow direction with reversal in sign of this difference, we use the definitions of Thual and McWilliams (1992)

$$q_o^+ = \frac{1}{2} (|q_o| + q_o)$$

$$q_o^- = \frac{1}{2} (|q_o| - q_o)$$

which automatically select the appropriate upstream differencing. Turbulent vertical diffusion is represented by entrainment velocities, respectively $w_i$ and $w_o$ inside (i) and outside (o) Saanich Inlet, and acting across the respective interfacial areas $A_i = A_1 = A_3$ and $A_o = A_2 = A_4$. $F S_2$ and $F S_4$, with units of salt per unit mass (salinity) per unit time, represent external forcing of the system through addition of freshwater and saltwater respectively to boxes 2 and 4.
Dividing eqns. (A1- A4) by $m_1$ and defining $\delta \equiv h/H = h_1/h_3 = h_2/h_4$ and $\Delta \equiv A_2/A_1$, the equations can be re-written as:

\[
\frac{dS_1}{dt} = q^+(S_3 - S_i) + q^-(S_2 - S_i) + W_i(S_3 - S_i) \tag{A5}
\]

\[
\frac{dS_2}{dt} = \frac{1}{\Delta} \left( q^+(S_1 - S_2) + q^-(S_4 - S_2) \right) + W_o(S_4 - S_2) - FS_2 \tag{A6}
\]

\[
\frac{dS_3}{dt} = \delta \left( q^+(S_4 - S_3) + q^-(S_1 - S_3) - W_i(S_3 - S_i) \right) \tag{A7}
\]

\[
\frac{dS_4}{dt} = \frac{\delta}{\Delta} \left( q^+(S_2 - S_4) + q^-(S_3 - S_4) \right) - \delta W_o(S_4 - S_2) + FS_4 \tag{A8}
\]

in which $W_{i/o} \equiv w_{i/o}/h$, and $q^{+/-} \equiv q^{+/-}/m_1$ have units of inverse time. Over an annual cycle, conservation of the total mass of salt in the system requires that $\frac{d}{dt} \left( \sum_{i=1}^{4} m_i S_i \right) = 0$, hence that $\int (m_i FS_i - m_j FS_j) = 0$, where the integral is over one year.

In keeping with the exploratory nature of the box model, we use a very simple functional form for the required (salinity) forcing function(s). Freshet is represented by a rectified cosine of temporal extent $T_w$, peaking at $T_p$, superimposed upon a much smaller, constant background "winter" level. The associated rate of change of salinity (in units of $[S]\text{s}^{-1}$) is

\[
FS_2(t) = F(t) = F_o \left[ 0.1 - 0.45(1 - \cos \left( \frac{2\pi(t - T_p)}{T_w} \right)) \right].
\]

Peak magnitude $F_o$ was estimated by converting average peak Fraser River flow of $Q_F = 6000 \text{ m}^3\text{s}^{-1}$ to equivalent salt forcing under the assumptions of Gill (1982): thus

\[
F_o = \frac{Q_F S_o}{V_o} = 5.3 \times 10^{-6} [S]\text{s}^{-1} \text{ where } V_o = 4 \times 10^{10} \text{ m}^3 \text{ is an estimate of the volume of the upper Strait of Georgia into which the Fraser River freshwater is distributed and } S_o = 35 \text{ is an average salinity of the Strait}. \text{ The required salt conservation is then simply enforced by setting } FS_4(t) = \delta F(t + \tau), \text{ where } \tau \text{ is a time offset. Fig.9a shows normalized forcing functions } F/F_o \text{ (plotting freshet as negative) for } T_p = \text{Year-day 181}, T_w = 60\text{d and } \tau = 60\text{d}.
References:


